

**WEEK 2: THE MATHEMATICS UNDERLYING GENERAL RELATIVITY****Recommended Reading:**

1. Roger D. Blandford and Kip S. Thorne, *Applications of Classical Physics* [cited henceforth as “Blandford and Thorne”], available on the web at <http://www.pma.caltech.edu/Courses/ph136/ph136.html>
  - a. Sections 23.1–23.3 of Chapter 23, “From Special to General Relativity” (version 0023.2)
  - b. Equation (24.39) and Section 24.6 of Chapter 24, “The Field Equations of General relativity” (version 0024.2).

Note: this introduction to the mathematics underlying general relativity assumes some prior familiarity with special relativity from a particular point of view, which is presented in Chapter 1 of Blandford and Thorne, “Physics in Flat Spacetime: Geometric Viewpoint” (available at above url). Some readers may wish to consult that chapter as they study Chapters 23 and 24; there are extensive cross references to it.

**Possible Supplementary Reading:**

2. Blandford and Thorne, Chapter 1, “Physics in Flat Spacetime: Geometric Viewpoint”. See note on item 1 above.
3. Bernard F. Schutz, *A First Course in General Relativity* (Cambridge University Press, 1985 & 1990), Chapters 2, 3, 5, 6. Not available on the web; a few copies are on sale at the bookstore, and one copy is on reserve in Millikan Library.
4. Charles W. Misner, Kip S. Thorne, and John A. Wheeler, *Gravitation* (Freeman, 1973), Chapters 2, 3, 5, 8. Not available on the web; a few copies are on sale at the bookstore, and one copy is on reserve in Millikan Library. Warning: This textbook is terribly out of date, as are *all* other advanced textbooks on general relativity. However, that does not affect the introduction to general relativity; only the applications are out of date (black holes, gravitational waves, cosmology, experimental tests, ...).

**Assignment, to be turned in at beginning of class on Wednesday 23 January by students registered in the course:**

- A. State what reading you have done, related to the course, during this past week.
- B. Work those exercises, from the list below, that are useful for you (i.e. that are at the appropriate level for you [neither much too hard nor too easy] and that have a ratio of grunge to learning that is reasonable.
- C. If A. and B. do not constitute enough to have taught you a reasonable amount about this week’s topic, then do one or more of the following:
  - i. If you already know a lot about this week’s topic, just say so and stop.
  - ii. Invent your own exercises and work them.
  - iii. Carry out further reading and state what you have done.

- iv. Seek private tutoring from a knowledgeable person about this week's topic.
- v. Pursue some other method of learning about this week's topic, and state what you have done.

## EXERCISES

Note: There are more exercises here than any single person is expected to work. Work only those exercises that are useful for you!

### Exercises filling in the gaps in Kip's Wednesday lecture

#### 1. Computation of components of a tensor

From the duality relation  $\vec{e}^\mu \cdot \vec{e}_\nu = \delta_\nu^\mu$  and expansions of a tensor  $\mathbf{T}(-, -, -)$  in terms of basis vectors, e.g.  $\mathbf{T} = T^{\alpha\beta}{}_\mu \vec{e}_\alpha \otimes \vec{e}_\beta \otimes \vec{e}^\mu$ , deduce that the components of a tensor can be computed by inserting basis vectors into its slots and lining up the indices, e.g.  $T^{\alpha\beta}{}_\mu = \mathbf{T}(\vec{e}^\alpha, \vec{e}^\beta, \vec{e}_\mu)$ .

#### 2. Raising and lowering of indices

From the properties of tensors discussed in exercise 1 and the definition of the metric in terms of the inner product,  $\mathbf{g}(\vec{A}, \vec{B}) = \vec{A} \cdot \vec{B}$ , show that indices on tensors can be raised and lowered using the metric components, e.g.  $T^{\alpha\beta}{}_\mu = g^{\alpha\rho} g_{\mu\sigma} T_\rho{}^{\beta\sigma}$ .

#### 3. Directional derivatives of bases

From the duality relation for bases and the definition  $\nabla_{\vec{e}_\alpha} \vec{e}_\beta \equiv \Gamma^\mu{}_{\beta\alpha} \vec{e}_\mu$  of the connection coefficients, show that  $\nabla_{\vec{e}_\alpha} \vec{e}^\rho \equiv -\Gamma^\rho{}_{\nu\alpha} \vec{e}^\nu$ .

#### 4. Connection coefficients for the orthonormal basis associated with circular polar coordinates, and their use

In Euclidean 2-space (a flat sheet of paper) construct circular polar coordinates  $(r, \phi)$ .

- a. Show that the basis

$$\mathbf{e}_{\hat{r}} \equiv \frac{\partial}{\partial r}, \quad \mathbf{e}_{\hat{\phi}} \equiv \frac{1}{r} \frac{\partial}{\partial \phi}$$

is orthonormal; i.e. in this basis the components of the metric are the Kronecker delta.

- b. By drawing pictures, deduce the values of all the connection coefficients for this basis.
- c. Let  $A^{\hat{\alpha}}$  be the components of a vector field  $\mathbf{A}$  in this basis. Using your connection coefficients, derive a formula for the divergence of  $\mathbf{A}$ ,  $\nabla \cdot \mathbf{A} = A^{\hat{\alpha}}{}_{;\hat{\alpha}}$  in terms of partial derivatives of the components of  $\mathbf{A}$ . Your answer should be the familiar formula

$$\nabla \cdot \mathbf{A} = \frac{1}{r} \frac{\partial(rA^{\hat{r}})}{\partial r} + \frac{1}{r} \frac{\partial A^{\hat{\phi}}}{\partial \phi}.$$

#### 5. Components of gradient of a tensor in terms of connection coefficients

By the same technique as is used in Eq. (23.29) of Blandford and Thorne, derive an expression for  $F^\alpha{}_{\beta;\mu}$  in terms of  $F^\alpha{}_{\beta,\mu}$  and the components of  $\mathbf{F}$  and the connection coefficients of the chosen basis.

**6. Components of commutator of two vector fields**

Let  $\vec{A}(\mathcal{P})$  and  $\vec{B}(\mathcal{P})$  be two vector fields. In an arbitrary basis (which might or might not be a coordinate basis), derive a formula for the components of the commutator  $[\vec{A}, \vec{B}]$  in terms of the components  $A^\alpha$  and  $B^\beta$ , their derivatives along the basis vectors,  $A^{\alpha, \mu}$  and  $B^{\beta, \nu}$ , and the basis's commutation coefficients  $c_{\alpha\beta\gamma}$ . In a coordinate basis your result should reduce to the one given by Kip in his lecture; cf. Eq. (23.24) of Blandford and Thorne.

**7. Formula for components of Riemann tensor in an arbitrary basis**

Exercise 24.8 of Blandford and Thorne

**8. Formula for components of Riemann tensor in a local Lorentz frame**

Show that in a local Lorentz frame in spacetime, the components of the Riemann tensor are given by Eq. (24.51) of Blandford and Thorne.

**Additional Exercises**

**8. Practice with frame-independent tensors**

Exercise 23.3 of Blandford and Thorne.

**9. Practice with index shuffling**

Let  $\mathbf{F}$  be a second-rank tensor that is antisymmetric under interchange of its slots, i.e.  $F_{\alpha\beta} = -F_{\beta\alpha}$ , and that satisfies the relation

$$F_{\alpha\beta;\gamma} + F_{\beta\gamma;\alpha} + F_{\gamma\alpha;\beta} = 0 .$$

Define  $J^\alpha$  by  $F^{\alpha\beta}{}_{;\beta} = 4\pi J^\alpha$ , and define

$$T^{\mu\nu} = \frac{1}{4\pi} \left( F^{\mu\alpha} F^\nu{}_\alpha - \frac{1}{4} g^{\mu\nu} F_{\alpha\beta} F^{\alpha\beta} \right) .$$

(Actually, the electric and magnetic fields can be embodied in an antisymmetric field tensor; if  $\mathbf{F}$  is that tensor, then  $\vec{J}$  is the charge-current 4-vector and  $\mathbf{T}$  is the electromagnetic stress-energy tensor.) Show that

$$T^{\alpha\beta}{}_{;\beta} = -F^{\alpha\beta} J_\beta .$$

Actually, this equation describes the the rate at which energy and momentum are transferred between the electromagnetic field and the charge-current distribution with which it interacts.